



k_{\perp} -factorization and heavy quark production

M. B. Gay Ducati

gay@if.ufrgs.br

GFPAE Group, Institute of Physics, Rio Grande do Sul U., Brazil

<http://www.if.ufrgs.br/gfpae>

Based on C. Brenner Mariotto, M. B. Gay Ducati, M. Machado, hep-ph/0208155

Overview:

- Motivation
- Heavy quark photoproduction in k_{\perp} -factorization approach
- Unintegrated gluon distributions
- Results
- Conclusions

Motivation

- Parton model (collinear approach)
 - NLO calculations
 - on-shell matrix elements, initial transverse momenta neglected

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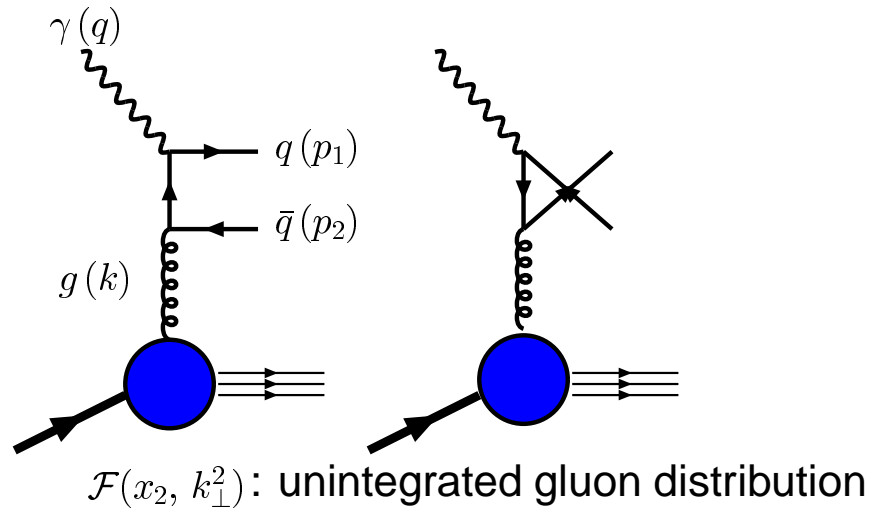
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- k_{\perp} -factorization approach (semihard interactions)
 - off-shell matrix elements
 - transverse momenta of incident partons included perturbatively
 - unintegrated parton distributions:
$$xg(x, \mu^2) \simeq \int_0^{\mu^2} dk_{\perp}^2 x \mathcal{A}(x, k_{\perp}^2, \mu^2):$$
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 $xg(x, \mu^2) \simeq \int_0^{\mu^2} dk_{\perp}^2 x \mathcal{A}(x, k_{\perp}^2, \mu^2)$: BFKL in asymptotic energy limit
 - NLO and some NNLO contributions on collinear approach already included in LO matrix elements
 - infrared sector and saturation effects

Heavy quark photoproduction in k_{\perp} -factorization approach



- Using Sudakov variables

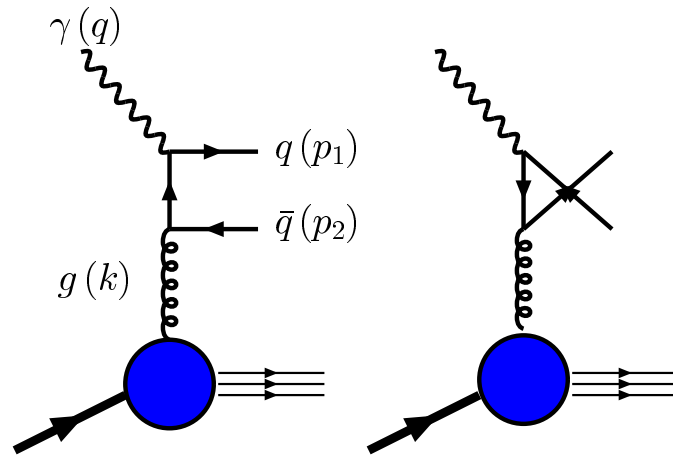
$$p_1 = \alpha_1 P_1 + \beta_1 P_2 + \mathbf{p}_{1\perp}$$

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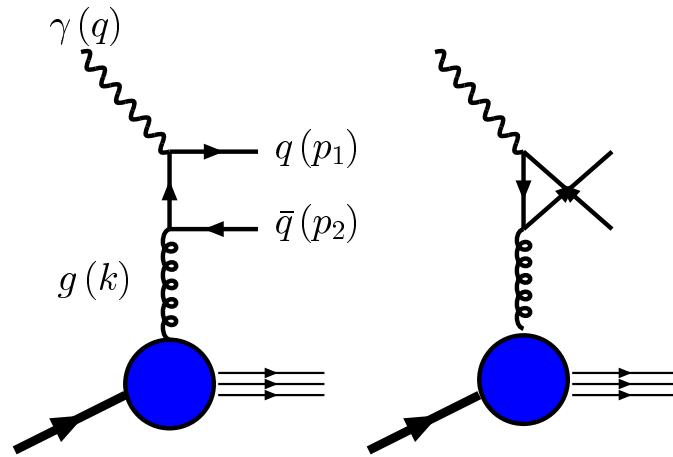
$$k = x_2 P_2 + \mathbf{k}_{\perp}$$

- in terms of transverse masses $m_{1,2\perp}^2 = m_Q^2 + \mathbf{p}_{1,2\perp}^2$ and heavy-quark rapidities $y_{1,2}^*$,

$$\alpha_1 = \frac{m_{1\perp}}{\sqrt{s}} \exp(y_1^*) \quad \beta_1 = \frac{m_{1\perp}}{\sqrt{s}} \exp(-y_1^*) \quad x_1 = \alpha_1 + \alpha_2$$

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- photoproduction: $q = P_1, x_1 = 1, \alpha_1 \equiv z, \alpha_2 \equiv 1 - z$
- $z, (1 - z)$: longitudinal momentum fraction carried by the heavy-quark having $\mathbf{p}_{1\perp}, \mathbf{p}_{2\perp}$

Heavy quark photoproduction in k_{\perp} -factorization approach

- direct interaction: (resolved component not included)

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where $D_1 \equiv p_{1\perp}^2 + m_Q^2$ and $D_2 \equiv (k_{\perp} - p_{1\perp})^2 + m_Q^2$.

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- input: $\mathcal{F}(x_2, k_{\perp}^2; \mu^2)$

Unintegrated gluon distributions

● GBW - Saturation Model

$$\mathcal{F}(x, \mathbf{k}_\perp^2) = \frac{3\sigma_0}{4\pi^2\alpha_s} R_0^2(x) \mathbf{k}_\perp^4 \exp(-R_0^2(x) \mathbf{k}_\perp^2)$$

$\sigma_0 = 29.12$ mb, $\alpha_s = 0.2$, $R_0(x) = \frac{1}{Q_0} \left(\frac{x}{x_0}\right)^{\lambda/2}$, parameterization including charm: $Q_0 = 1$ GeV, $\lambda = 0.277$ and $x_0 = 0.41 \cdot 10^{-4}$

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- **Derivative of integrated gluon density** (*d-Gluon*)

$$\mathcal{G}(x, \mathbf{k}_\perp^2) =$$

$$+ \frac{\partial [x G(x, \mathbf{k}_\perp^2)]}{\partial \ln \mathbf{k}_\perp^2} \Theta(\mathbf{k}_\perp^2 - Q_0^2) \quad (\text{pQCD})$$

$$xG(x, \mu^2): \text{GRV98 LO}, Q_0^2 = 0.8 \text{ GeV}^2$$

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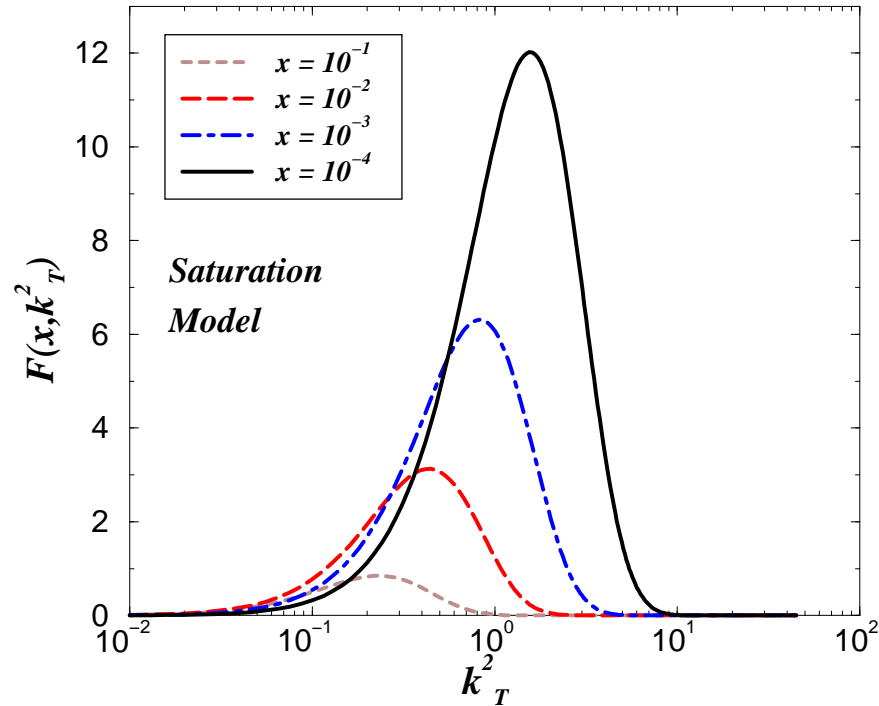
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$$\mathcal{G}(x, \mathbf{k}_\perp^2) = \mathbf{k}_\perp^2 \left. \frac{\partial [x G(x, \mathbf{k}_\perp^2)]}{\partial \mathbf{k}_\perp^2} \right|_{\mathbf{k}_\perp^2 = Q_0^2} \Theta(Q_0^2 - \mathbf{k}_\perp^2) \quad (\text{prescription})$$
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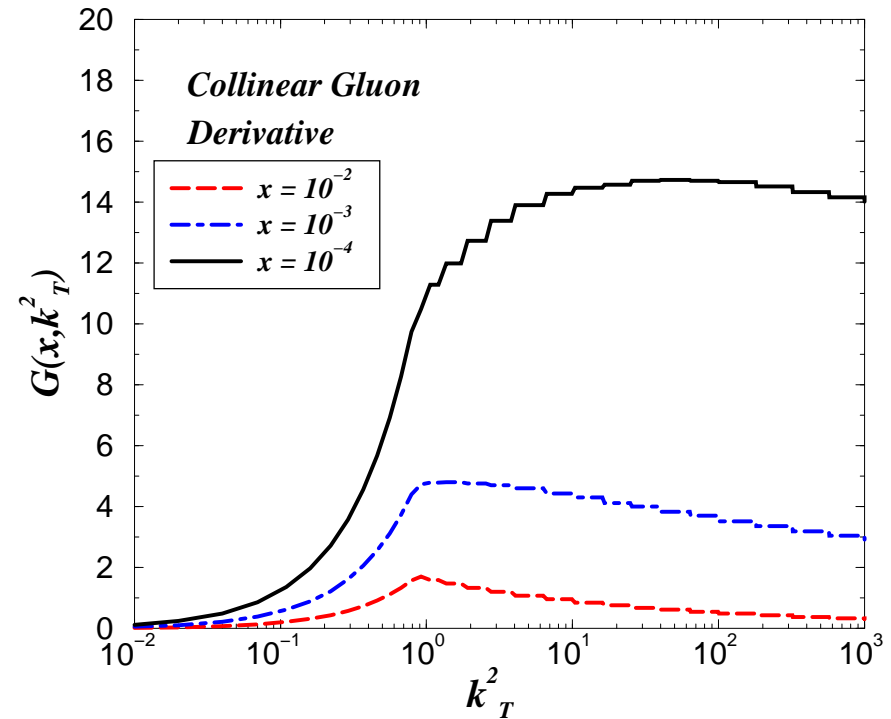
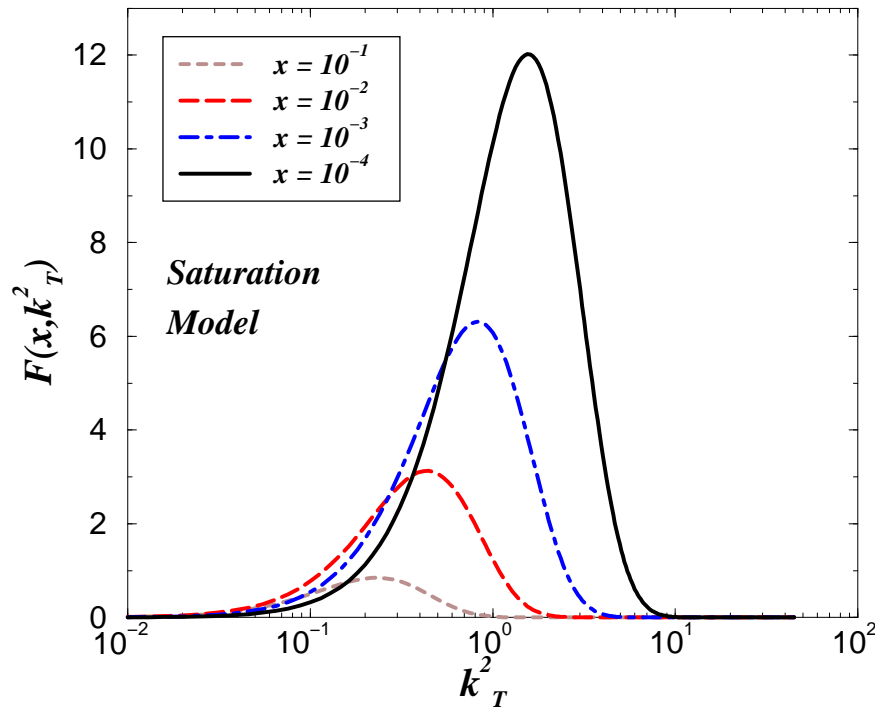
$xG(x, \mu^2)$: GRV98 LO, $Q_0^2 = 0.8 \text{ GeV}^2$

Unintegrated gluon distributions:



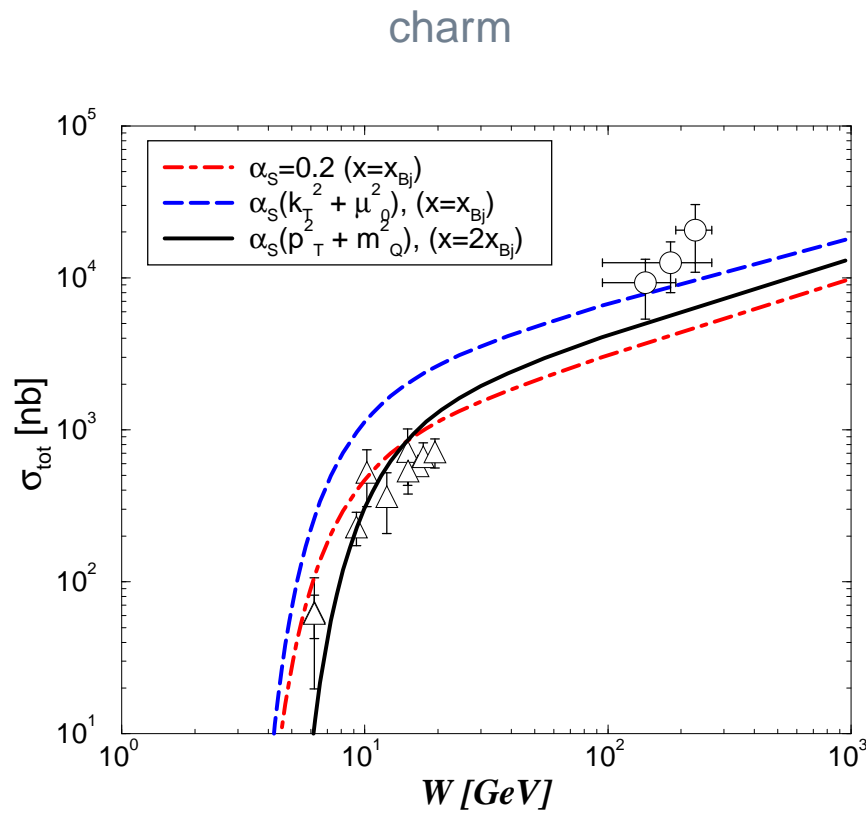
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Unintegrated gluon distributions:



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- d-Gluon broader in k_{\perp}^2

Total cross section, GBW

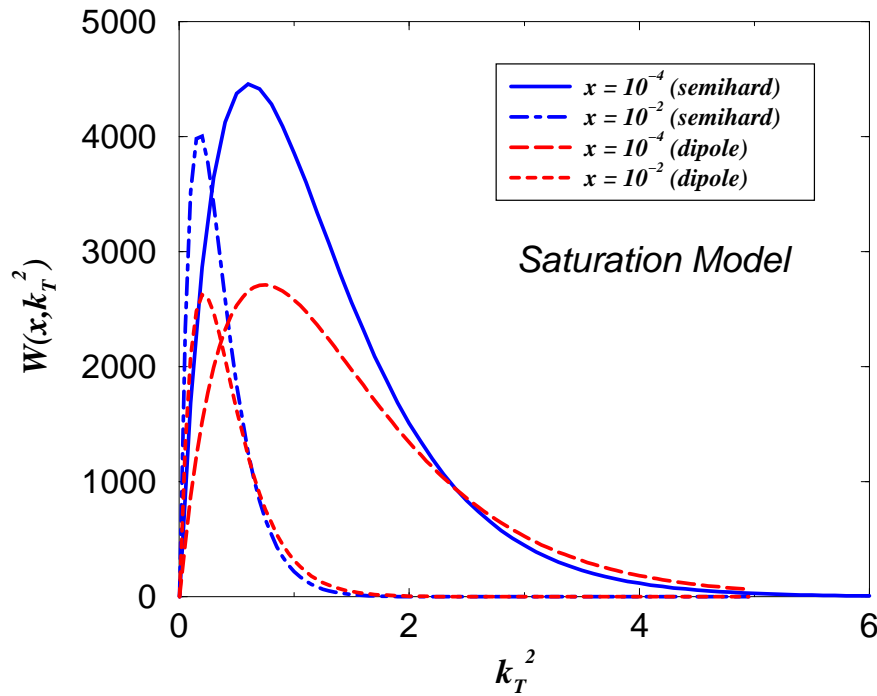


Sensitivity to energy scale and gluon momentum fraction

- $\alpha_s = 0.2, x_2 = x_{Bj}$ reproduces results of dipole approach (Szczurek)
- More conservative choice $\mu^2 = p_{\perp}^2 + m_Q^2, x_2 = 2 x_{Bj}$

$W(x, k_{\perp}^2) \rightarrow k_{\perp}$ profile, charm

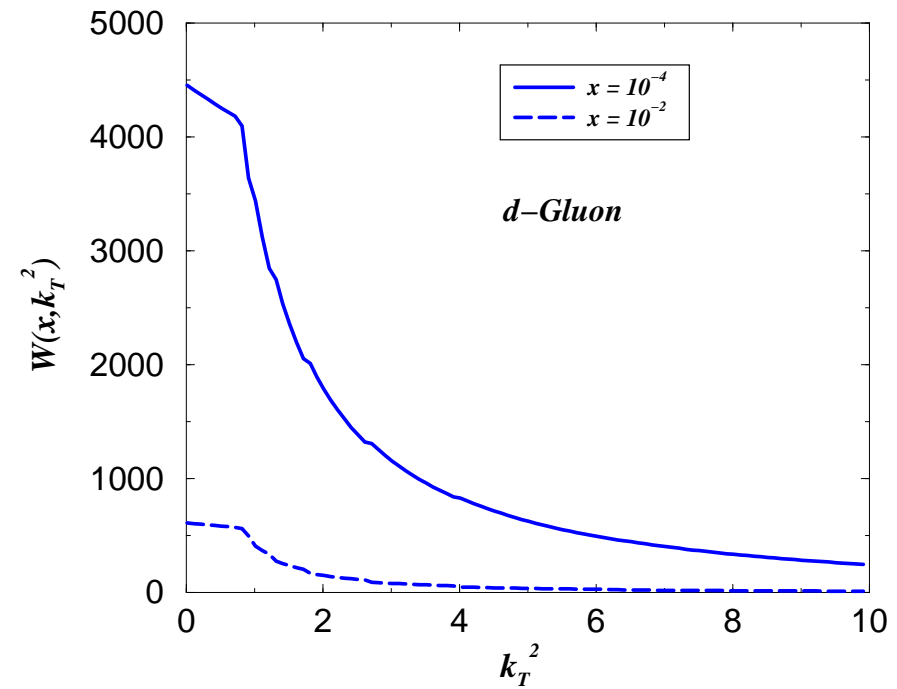
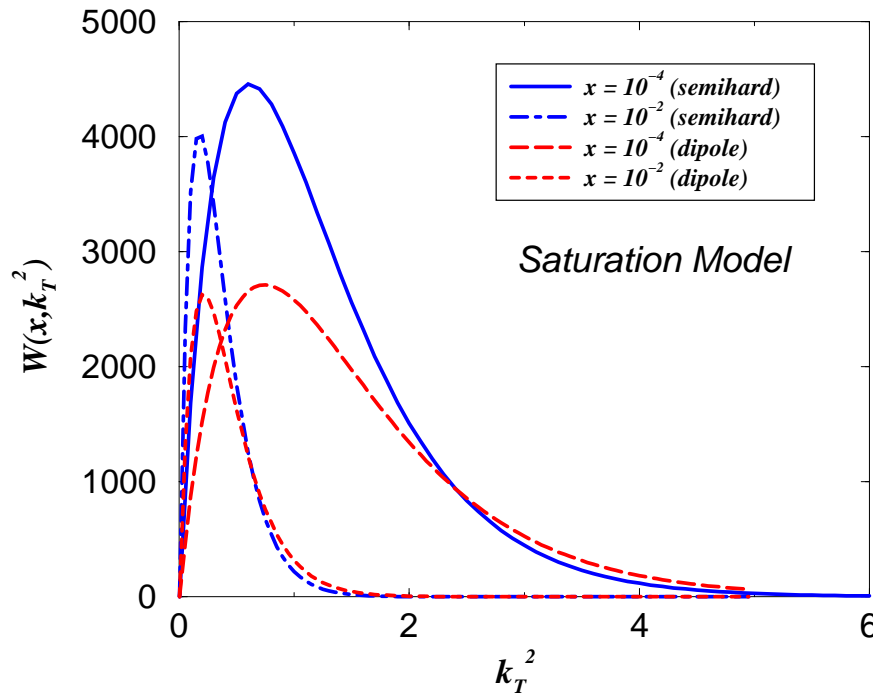
fixed target energies ($x = 10^{-2}$), high energies ($x = 10^{-4}$)



- Saturation Model: Dipole approximation (fixed α_s and $x_2 = x_{Bj}$)
Semihard approach ($\alpha_s(\mathbf{p}_{\perp}^2 + m_Q^2)$, $x_2 = 2x_{Bj}$): larger normalization

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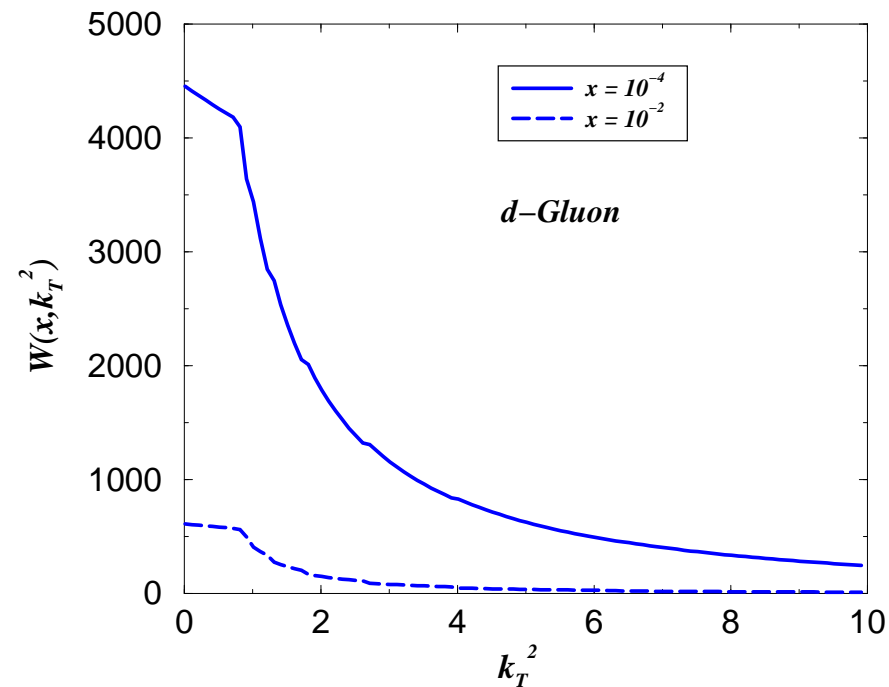
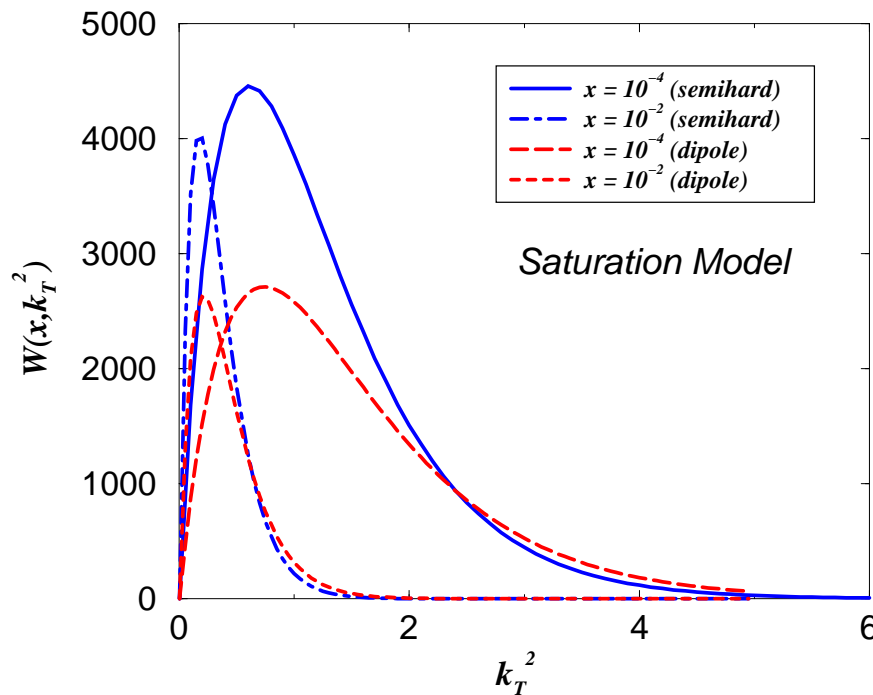
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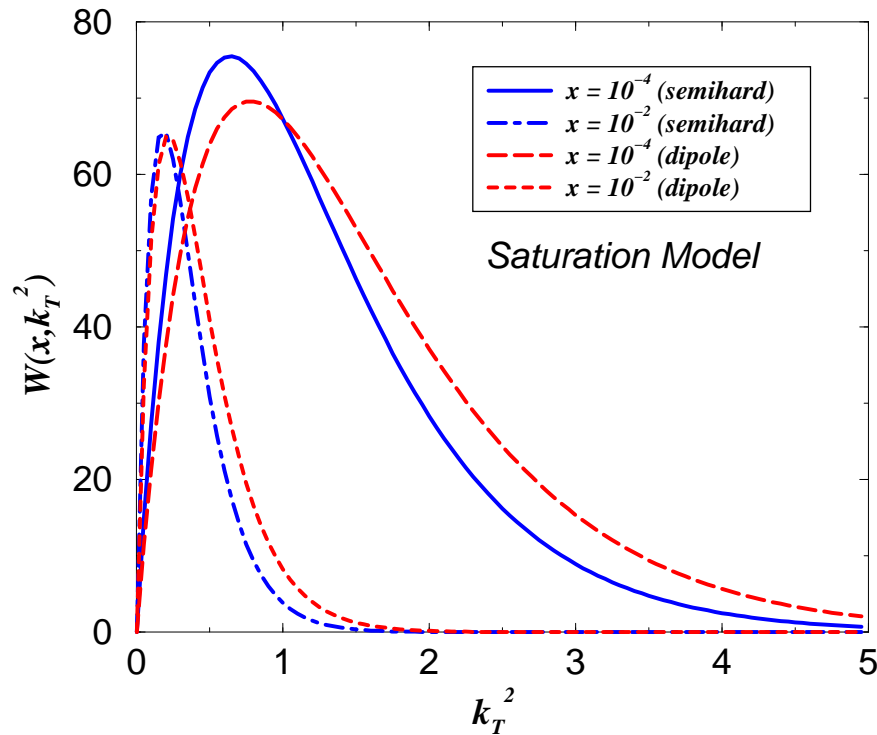
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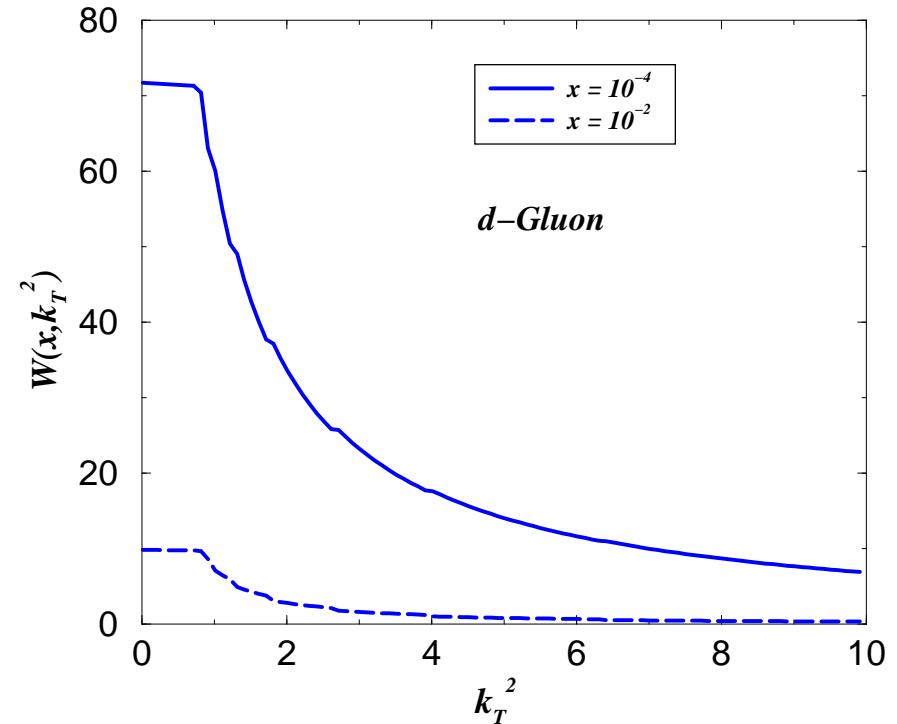
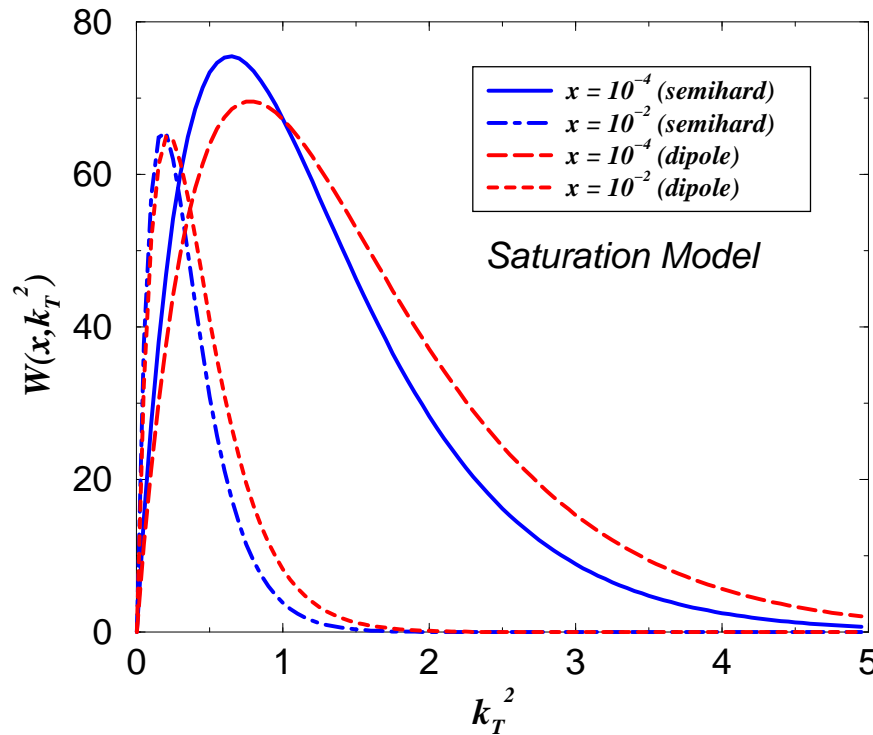
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- dominance of small k_{\perp}^2 region, infrared region under control

$W(x, k_{\perp}^2) \rightarrow k_{\perp}$ profile, bottom:



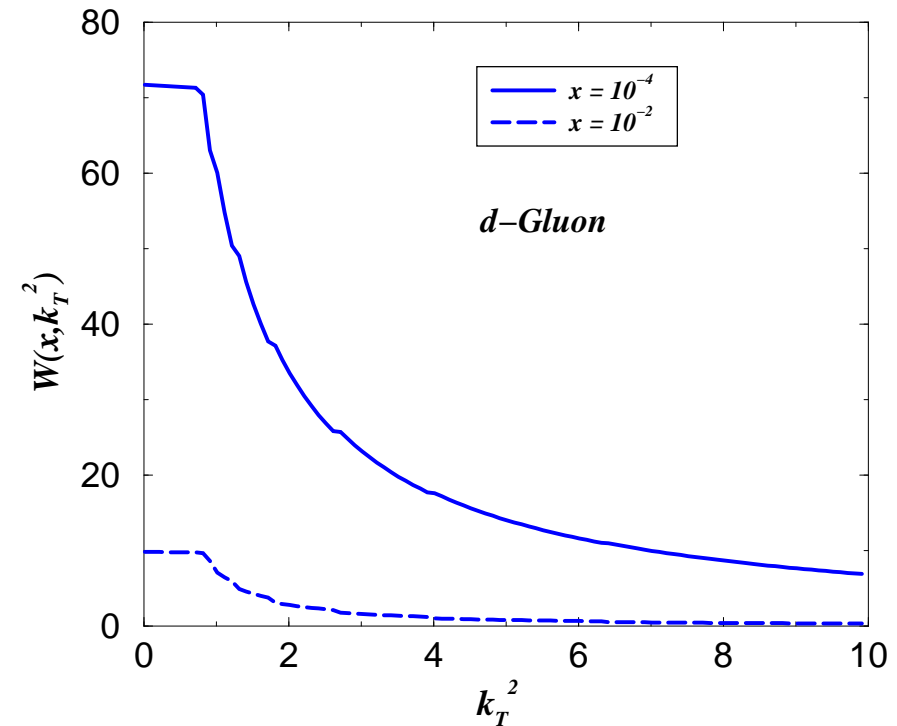
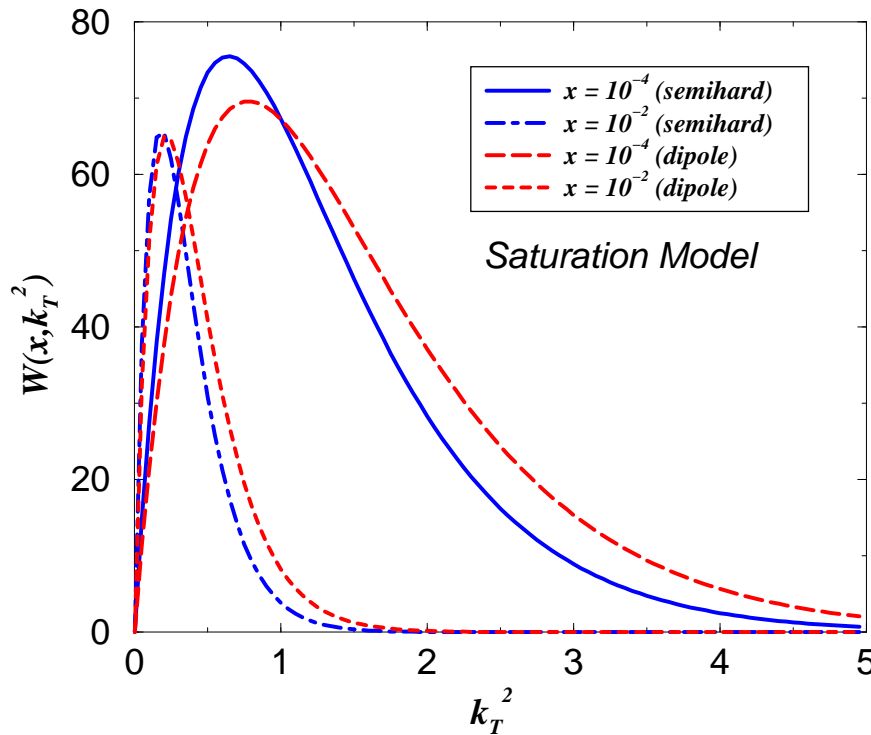
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(large m_b at the scale $\mu^2 = p_{\perp}^2 + m_Q^2$)

$W(x, k_{\perp}^2) \rightarrow k_{\perp}$ profile, bottom:



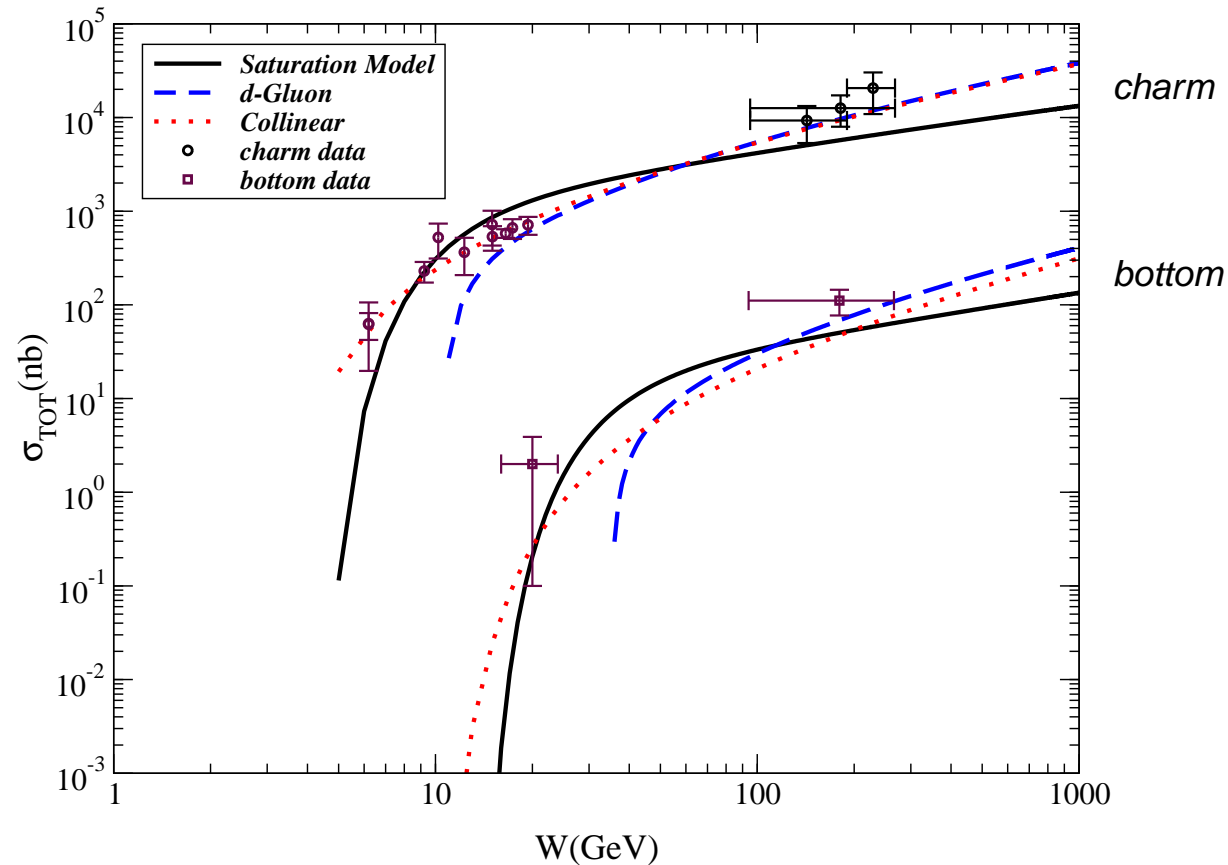
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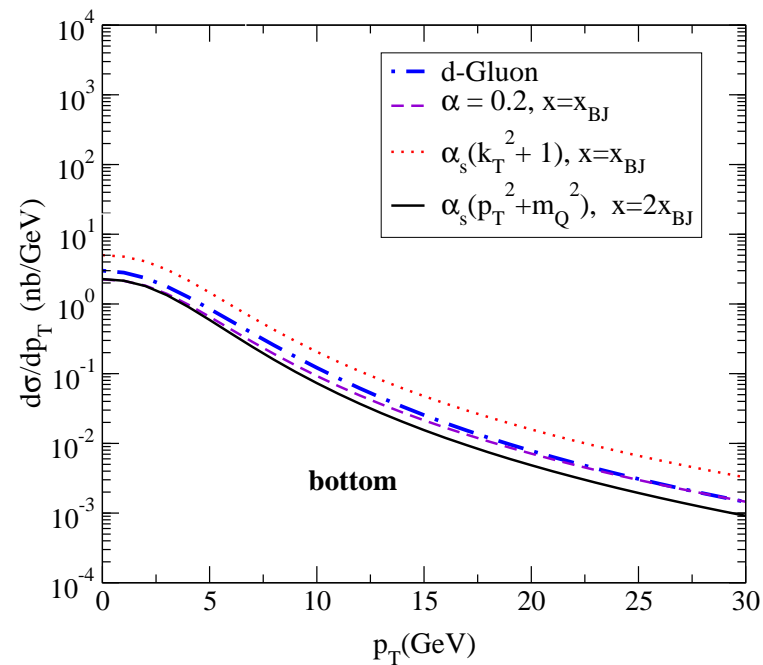
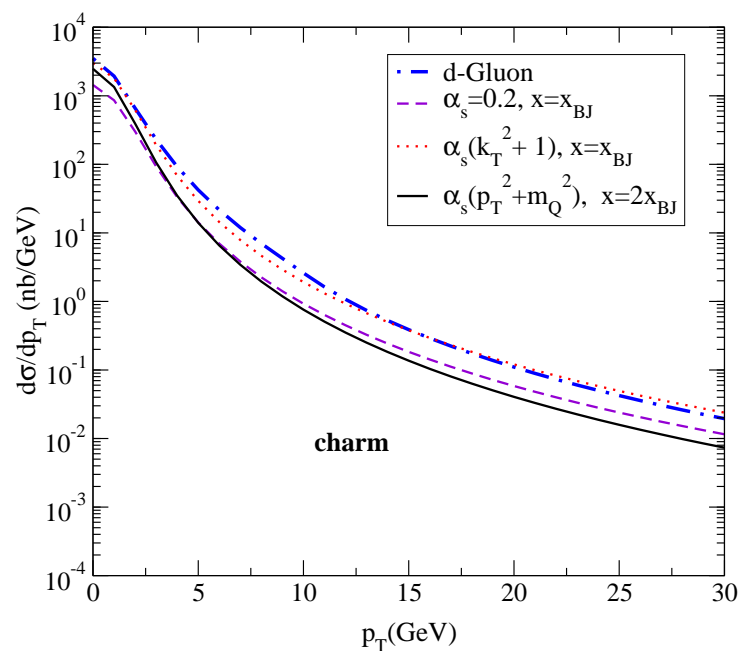
σ_{TOT} , Saturation Model x d-Gluon:



- Saturation Model: somewhat below high energy data (lack of QCD evolution)
- d-Gluon: better description of high energy data
lower energies: lack of non-singlet (valence) content, $\rightarrow dG < 0$

p_T distributions

$W = 200 \text{ GeV}$, saturation model X d-Gluon



- usual fall off at large transverse momentum
- finite and well-controlled behavior at $p_T = 0$

Conclusions

- k_T factorization: NLO and NNLO collinear diagrams in LO
- Several parameterizations for \mathcal{F} in the market:
 - Derivative of integrated gluon density**: DGLAP
 - GBW**: consistent with small- x DIS and diffraction
- smooth connection between soft and hard regimes
- profile function: important below saturation scale
- fixed $\alpha_s = 0.2$ and $\bar{x} = x_{Bj}$ underestimates high energy data (consistent with dipole approach)
- more conservative choice $\alpha_s(p_T^2 + m_Q^2)$, $\bar{x} = 2x_{Bj}$ improves slightly the description
- p_T spectrum finite at $p_T \rightarrow 0$